

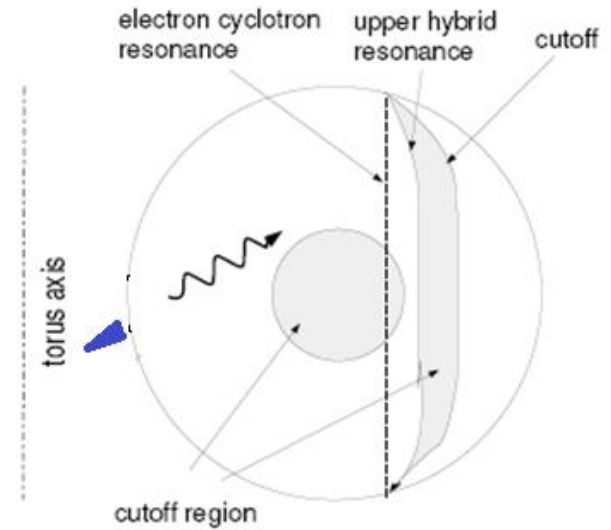
# Modelling of Electron Cyclotron Resonance Heating for COMPASS-U

**Michal Farník**

FTTF workshop  
Mariánská 2019

- **Motivation**
- **Electron Cyclotron Wave Physics**
- **ECRH for COMPASS-U**
- **SPECE code – Ray-Tracing**
- **Preliminary Results**
- **Summary**

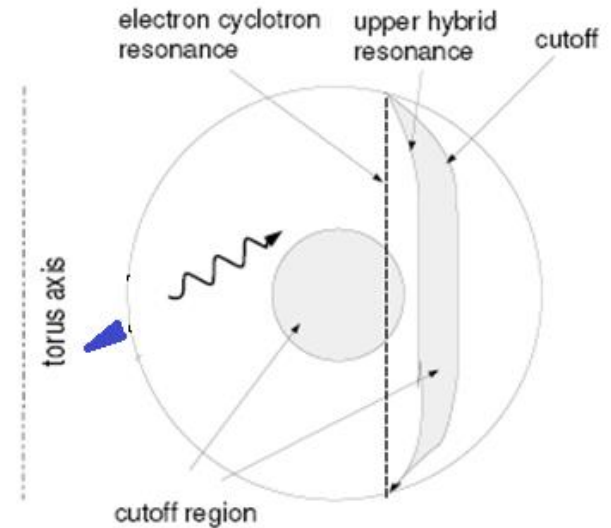
**COMPASS-U will be equipped with an NBI heating system.  
An additional ECRH heating is planned on the tokamak.**



**COMPASS-U will be equipped with an NBI heating system.  
An additional ECRH heating is planned on the tokamak.**

## Main features of the ECRH:

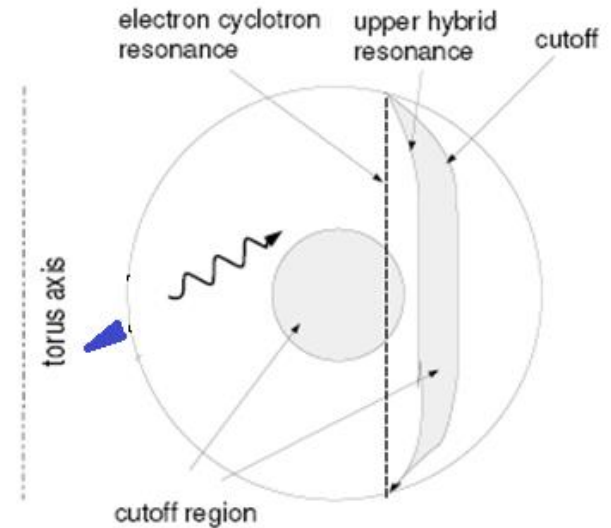
- Heating



**COMPASS-U will be equipped with an NBI heating system.  
An additional ECRH heating is planned on the tokamak.**

## Main features of the ECRH:

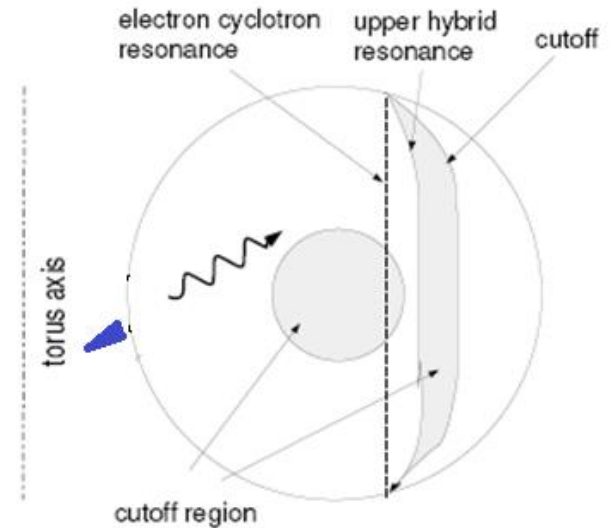
- Heating
- Improving H-mode performance



**COMPASS-U will be equipped with an NBI heating system.  
An additional ECRH heating is planned on the tokamak.**

## Main features of the ECRH:

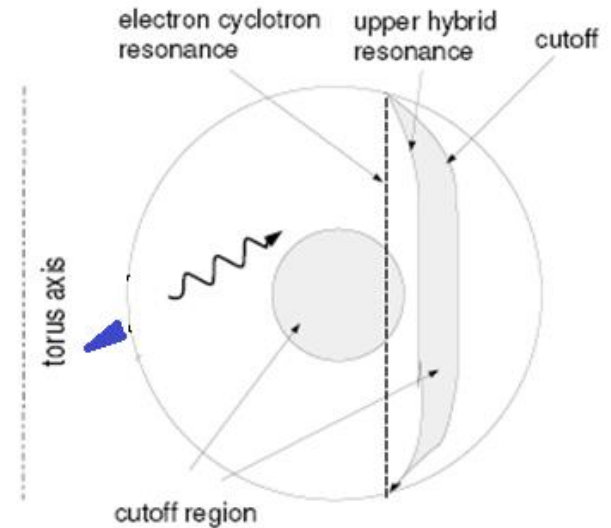
- Heating
- Improving H-mode performance
- Preventing impurity accumulation in the core



**COMPASS-U will be equipped with an NBI heating system.  
An additional ECRH heating is planned on the tokamak.**

## Main features of the ECRH:

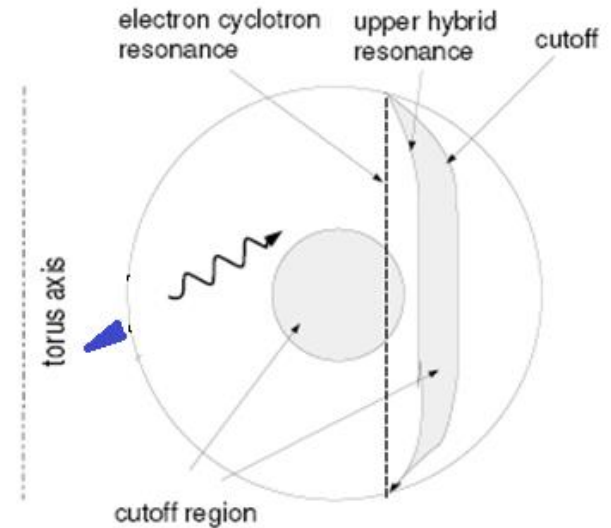
- Heating
- Improving H-mode performance
- Preventing impurity accumulation in the core
- Suppressing Neoclassical Tearing Modes (NTMs)



**COMPASS-U will be equipped with an NBI heating system.  
An additional ECRH heating is planned on the tokamak.**

## Main features of the ECRH:

- Heating
- Improving H-mode performance
- Preventing impurity accumulation in the core
- Suppressing Neoclassical Tearing Modes (NTMs)
- Current Drive (ECCD) – profile tailoring





$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

$$\mathbf{j} = -n_e e \mathbf{v} = \boldsymbol{\sigma} \cdot \mathbf{E}$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

$$\mathbf{j} = -n_e e \mathbf{v} = \boldsymbol{\sigma} \cdot \mathbf{E}$$

$$-i\omega m_e \mathbf{v}_e = -e(\mathbf{E} + \mathbf{v}_e \times \mathbf{B}_0)$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

$$\mathbf{j} = -n_e e \mathbf{v} = \boldsymbol{\sigma} \cdot \mathbf{E}$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$-i\omega m_e \mathbf{v}_e = -e (\mathbf{E} + \mathbf{v}_e \times \mathbf{B}_0)$$

$$\det(\mathbb{D}) = D(\mathbf{k}, \omega) \equiv \frac{c^4}{\omega^4} k_{\perp}^4 \epsilon_{11} - \frac{c^2}{\omega^2} k_{\perp}^2 \left[ \epsilon_{12}^2 + (\epsilon_{11} + \epsilon_{33}) \left( \epsilon_{11} - \frac{c^2}{\omega^2} k_{\parallel}^2 \right) \right] + \epsilon_{33} \left[ \left( \epsilon_{11} - \frac{c^2}{\omega^2} k_{\parallel}^2 \right)^2 + \epsilon_{12}^2 \right] = 0.$$

$$\nabla \times \mathbf{B} - \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} = \mu_0 \mathbf{j}$$

$$\mathbf{j} = -n_e e \mathbf{v} = \boldsymbol{\sigma} \cdot \mathbf{E}$$

$$\nabla \times \mathbf{E} + \frac{\partial \mathbf{B}}{\partial t} = 0$$

$$-i\omega m_e \mathbf{v}_e = -e (\mathbf{E} + \mathbf{v}_e \times \mathbf{B}_0)$$

$$\det(\mathbb{D}) = D(\mathbf{k}, \omega) \equiv \frac{c^4}{\omega^4} k_{\perp}^4 \epsilon_{11} - \frac{c^2}{\omega^2} k_{\perp}^2 \left[ \epsilon_{12}^2 + (\epsilon_{11} + \epsilon_{33}) \left( \epsilon_{11} - \frac{c^2}{\omega^2} k_{\parallel}^2 \right) \right] + \epsilon_{33} \left[ \left( \epsilon_{11} - \frac{c^2}{\omega^2} k_{\parallel}^2 \right)^2 + \epsilon_{12}^2 \right] = 0.$$

### APPLETON-HARTREE FORMULA

$$N^2 = 1 - \frac{2X(1-X)}{2(1-X) - Y^2 \sin^2 \theta \pm \left[ Y^4 \sin^4 \theta + 4Y^2(1-X)^2 \cos^2 \theta \right]^{1/2}}$$

$$X = (\omega_p / \omega)^2 \text{ a } Y = \omega_c / \omega$$

## Perpendicular propagation

Plasma cutoff  $\omega = \omega_O = \omega_{pe},$

Right – hand cutoff  $\omega = \omega_R = \left( \frac{\omega_{ce}^2}{4} + \omega_{pe}^2 \right)^{1/2} + \frac{\omega_{ce}}{2},$

Left – hand cutoff  $\omega = \omega_L = \left( \frac{\omega_{ce}^2}{4} + \omega_{pe}^2 \right)^{1/2} - \frac{\omega_{ce}}{2}.$

Upper hybrid resonance  $\omega = \omega_{UH} = (\omega_{ce} + \omega_{pe})^{1/2}$

## Perpendicular propagation

Plasma cutoff  $\omega = \omega_O = \omega_{pe},$

Right – hand cutoff  $\omega = \omega_R = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} + \frac{\omega_{ce}^2 - \omega_{pe}^2}{2} \right)^{1/2}$

Left – hand cutoff  $\omega = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} - \frac{\omega_{ce}^2 - \omega_{pe}^2}{2} \right)^{1/2} = \frac{\omega_{ce}}{2}.$

Upper hybrid resonance  $\omega = \omega_{UH} = (\omega_{ce} + \omega_{pe})^{1/2}$

**NO CYCLOTRON RESONANCE?!**



## Perpendicular propagation

Plasma cutoff  $\omega = \omega_O = \omega_{pe},$

Right – hand cutoff  $\omega = \omega_R = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} \right)^{1/2}$

Left – hand cutoff  $\omega = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} \right)^{1/2} - \frac{\omega_{ce}}{2}.$

**NO CYCLOTRON RESONANCE?!**

Upper hybrid resonance  $\omega = \omega_{UH} = (\omega_{ce} + \omega_{pe})^{1/2}$

## Hot plasma – Kinetic theory

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - e(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{p}} = 0$$

## Perpendicular propagation

Plasma cutoff  $\omega = \omega_O = \omega_{pe},$

Right – hand cutoff  $\omega = \omega_R = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} \right)^{1/2}$

Left – hand cutoff  $\omega = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} \right)^{1/2} - \frac{\omega_{ce}}{2}.$

**NO CYCLOTRON RESONANCE?!**

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - e(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{p}} = 0$$

## Hot plasma – Kinetic theory

### RESONANCE CONDITION

$$\omega = \frac{n \cdot \omega_c(B, v)}{1 - \beta_{\parallel} N \cos(\theta)} + k_{\parallel} v_{\parallel}, \quad \omega_c(B, v) = \frac{eB}{m_e(v)}$$

## Perpendicular propagation

Plasma cutoff  $\omega = \omega_O = \omega_{pe}$ ,

Right – hand cutoff  $\omega = \omega_R = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} + \frac{\omega_{ce}^2 - \omega_{pe}^2}{2} \right)^{1/2}$  Upper hybrid resonance  $\omega = \omega_{UH} = (\omega_{ce} + \omega_{pe})^{1/2}$

Left – hand cutoff  $\omega = \left( \frac{\omega_{ce}^2 + \omega_{pe}^2}{2} - \frac{\omega_{ce}^2 - \omega_{pe}^2}{2} \right)^{1/2} - \frac{\omega_{ce}}{2}$ .

**NO CYCLOTRON RESONANCE?!**

## Hot plasma – Kinetic theory

$$\frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - e(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f}{\partial \mathbf{p}} = 0$$

### RESONANCE CONDITION

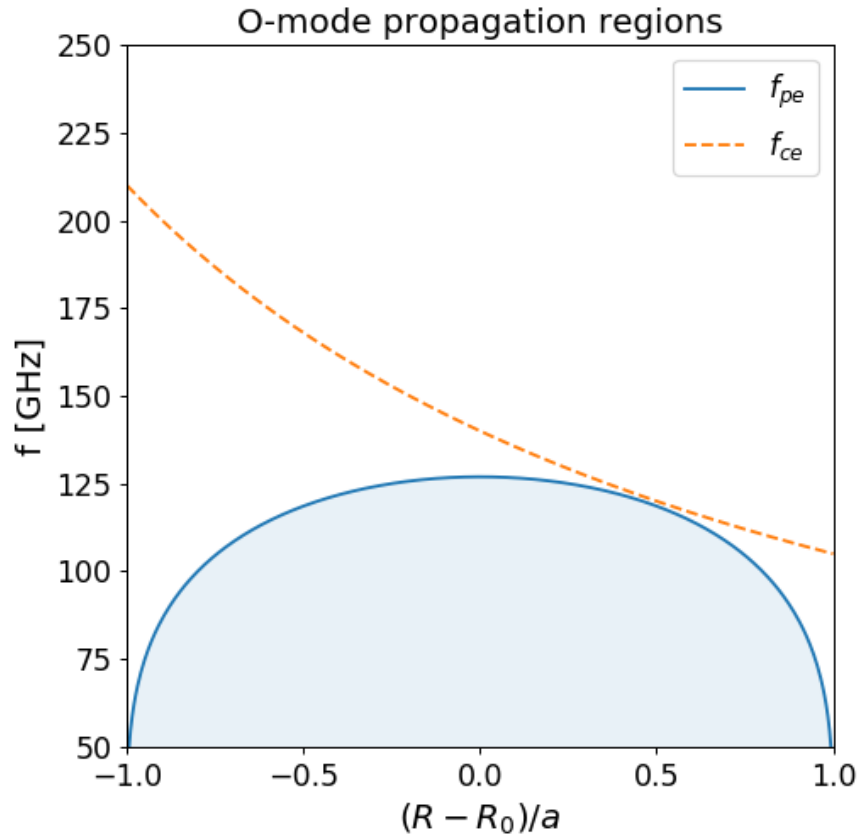
$$\omega = \frac{n \cdot \omega_c(B, v)}{1 - \beta_{\parallel} N \cos(\theta)} + k_{\parallel} v_{\parallel}, \quad \omega_c(B, v) = \frac{eB}{m_e(v)}$$

***n*-th harmonic**

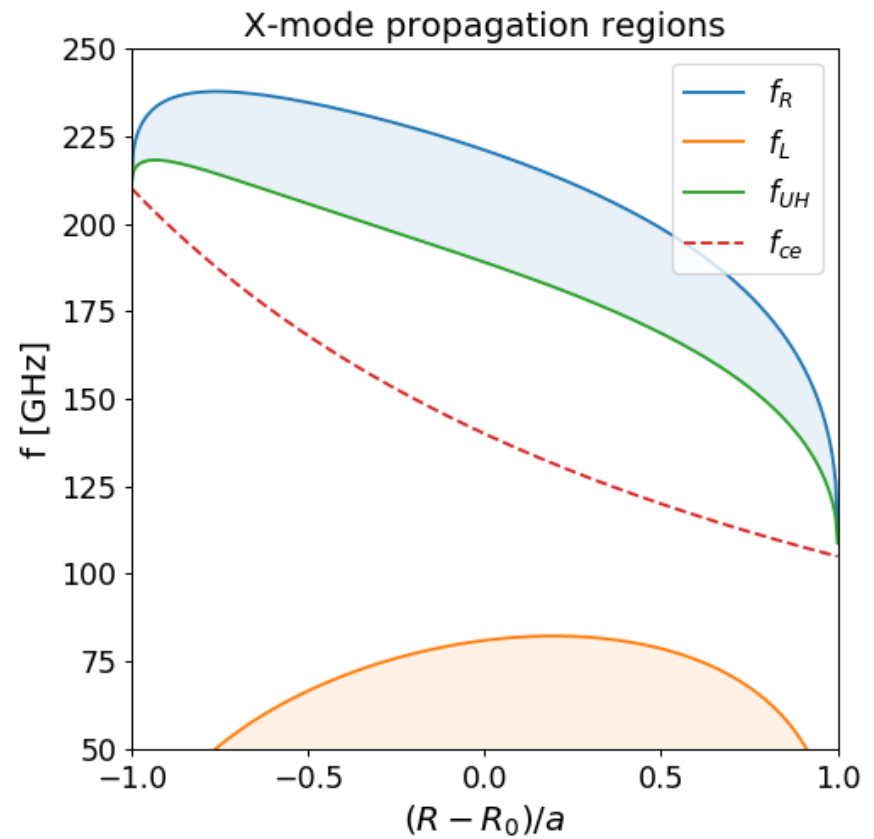
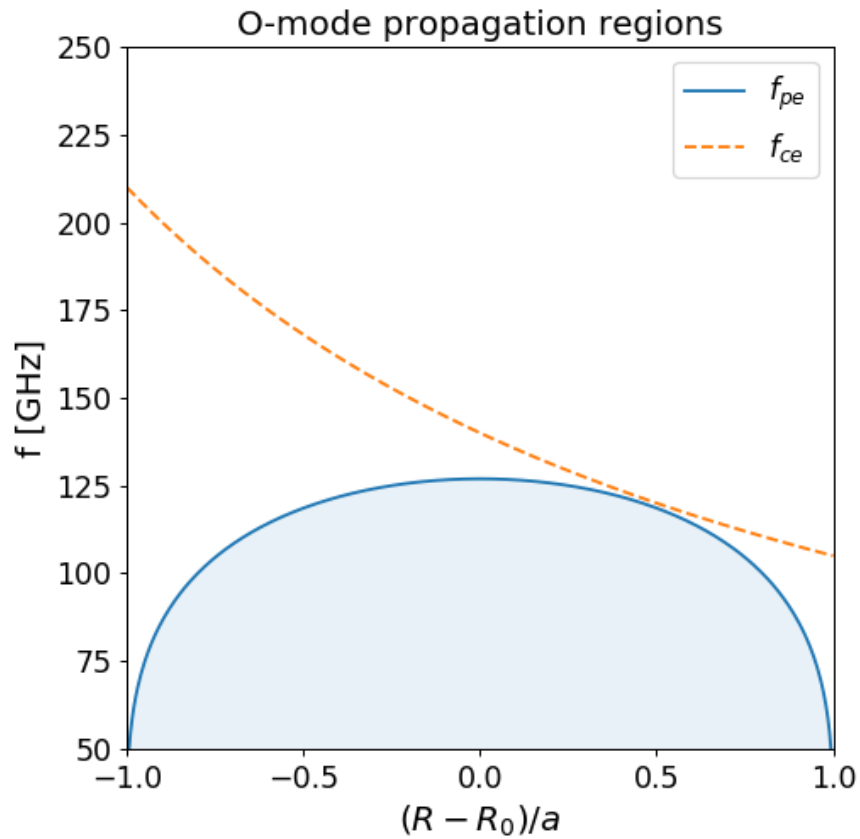
**Doppler shift**

**Relativistic mass**

## COMPASS-U



# COMPASS-U

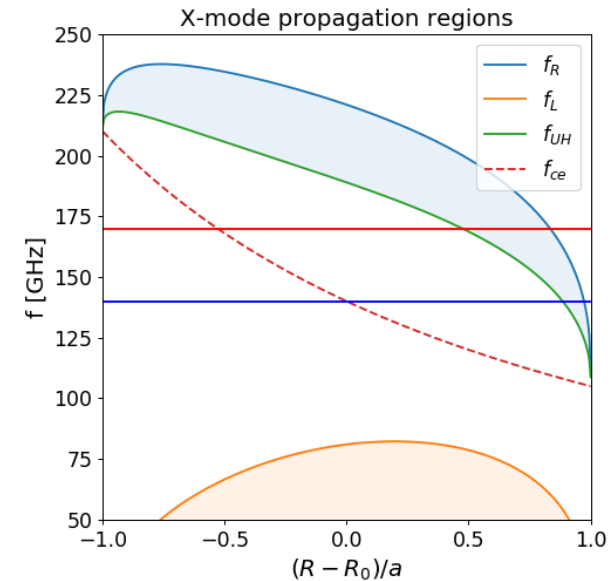
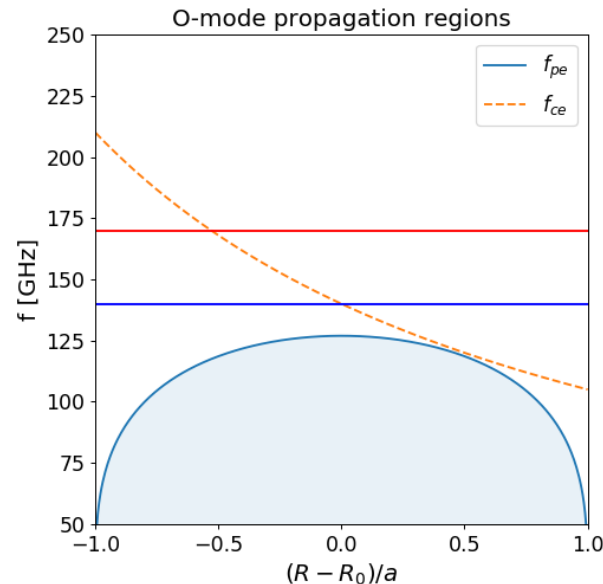


- **COMPASS-U**: DEMO relevant research ( $R_0 = 0.84$  m,  $a = 0.28$  m)  
 $B_T$ : **5 T**,  $I_p$ : **2 MA**, discharge length 1-5 s,  $n_e$ :  **$\sim 10^{20}$  m<sup>-3</sup>**  
 $P_{\text{NBI}} = 4\text{-}5$  MW,  **$P_{\text{ECRH}} = 4$  MW**

- **COMPASS-U**: DEMO relevant research ( $R_0 = 0.84$  m,  $a = 0.28$  m)  
 $B_T$ : **5 T**,  $I_p$ : **2 MA**, discharge length 1-5 s,  $n_e$ :  $\sim 10^{20}$  m<sup>-3</sup>  
 $P_{\text{NBI}} = 4\text{-}5$  MW,  $P_{\text{ECRH}} = 4$  MW
- **1<sup>st</sup> harmonic** heating (high  $B_T$ )
- Considered **140GHz** (ASDEX-U) or **170GHz** (ITER) systems

- COMPASS-U**: DEMO relevant research ( $R_0 = 0.84$  m,  $a = 0.28$  m)  
 $B_T$ : 5 T,  $I_p$ : 2 MA, discharge length 1-5 s,  $n_e$ :  $\sim 10^{20}$  m $^{-3}$   
 $P_{\text{NBI}} = 4\text{-}5$  MW,  $P_{\text{ECRH}} = 4$  MW
- 1<sup>st</sup> harmonic heating (high  $B_T$ )
- Considered 140GHz (ASDEX-U) or 170GHz (ITER) systems
- O-mode: cutoff density  $\approx 2.5 \cdot 10^{20}$  m $^{-3}$  for 140 GHz  
 X-mode: core achievable from the HFS

**Reminder:**





## Difficulties

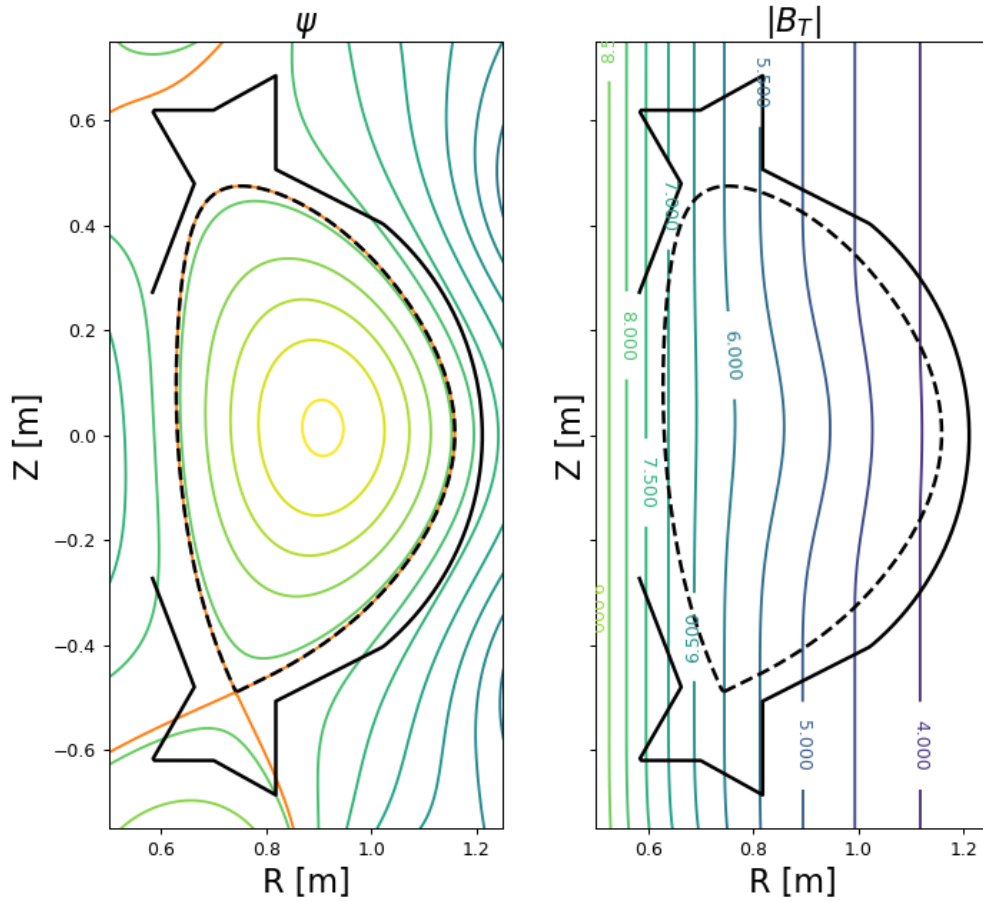
- **High densities** – above the cutoff condition in H-mode

## Difficulties

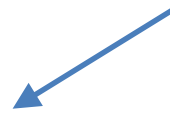
- **High densities** – above the cutoff condition in H-mode
- **Higher harmonic heating not possible** with current gyrotron technology (250 GHz gyrotron **GYCOM** in testing phase)

## Difficulties

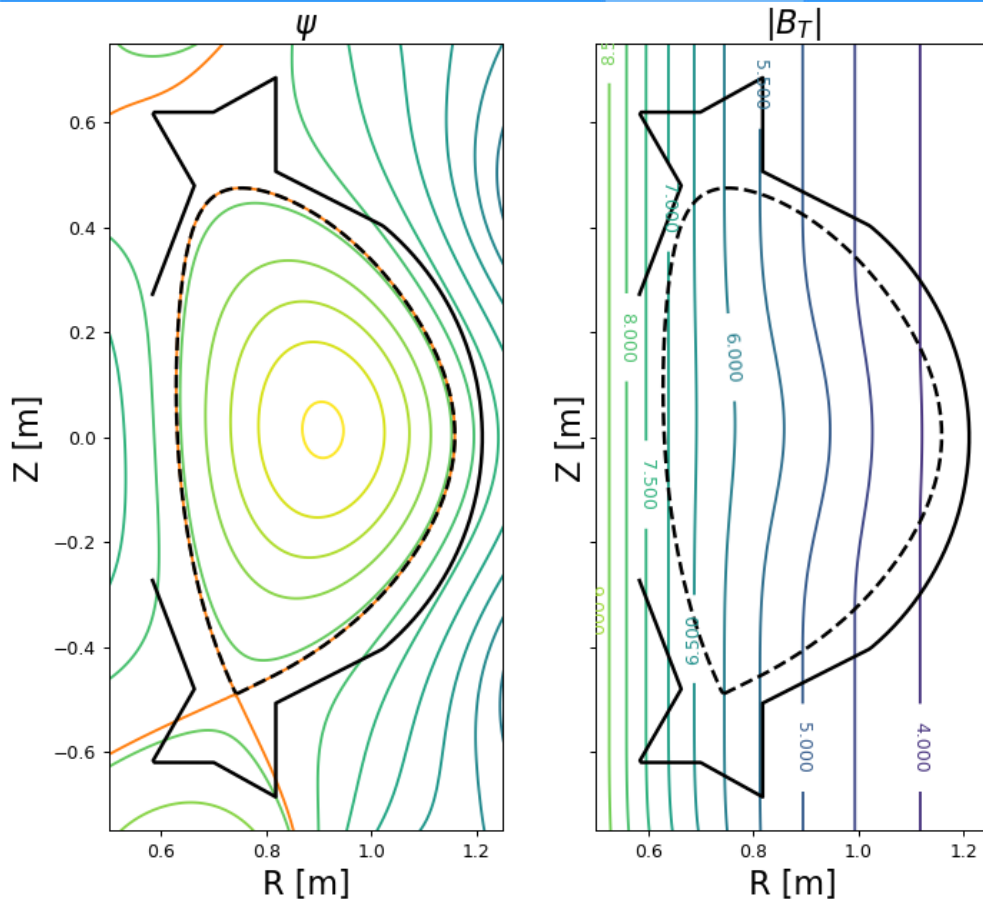
- **High densities** – above the cutoff condition in H-mode
- **Higher harmonic heating not possible** with current gyrotron technology (250 GHz gyrotron **GYCOM** in testing phase)
- **1<sup>st</sup> harmonic X-mode** is optically thin thus **not suitable** for the ECRH
- HFS launcher or mirror technically difficult



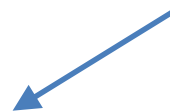
**Baseline 5T scenario equilibrium**



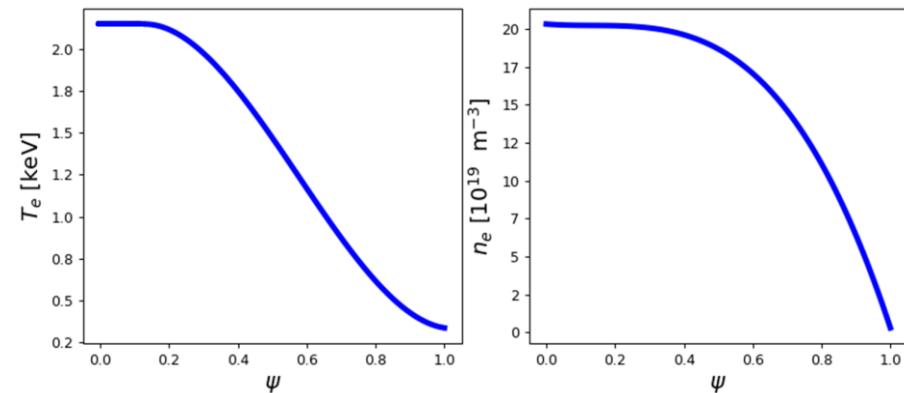
Credit: PLEQUE – L. Kripner



**Baseline 5T scenario equilibrium**

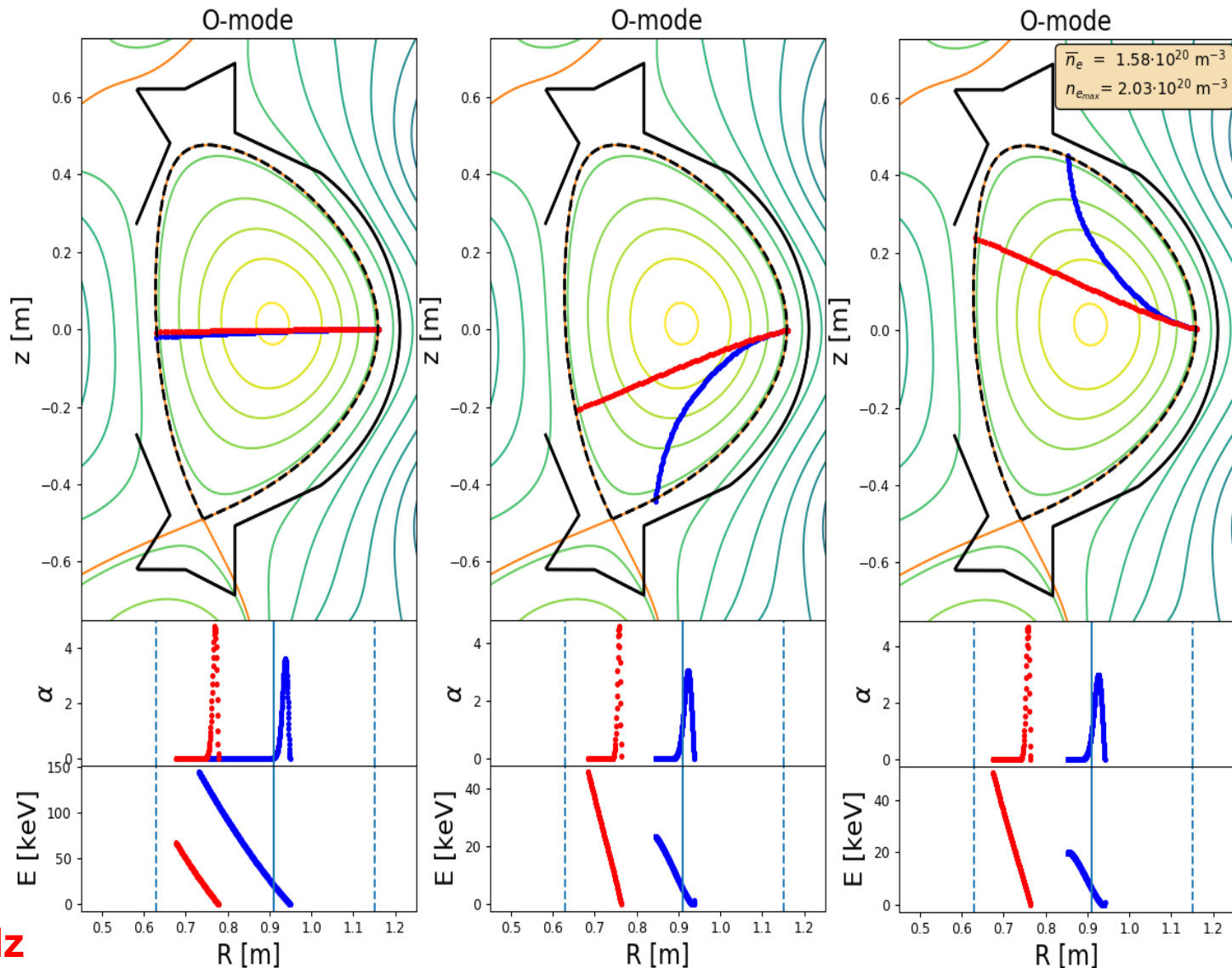


**Temperature and density profiles**



Credit: PLEQUE – L. Kripner

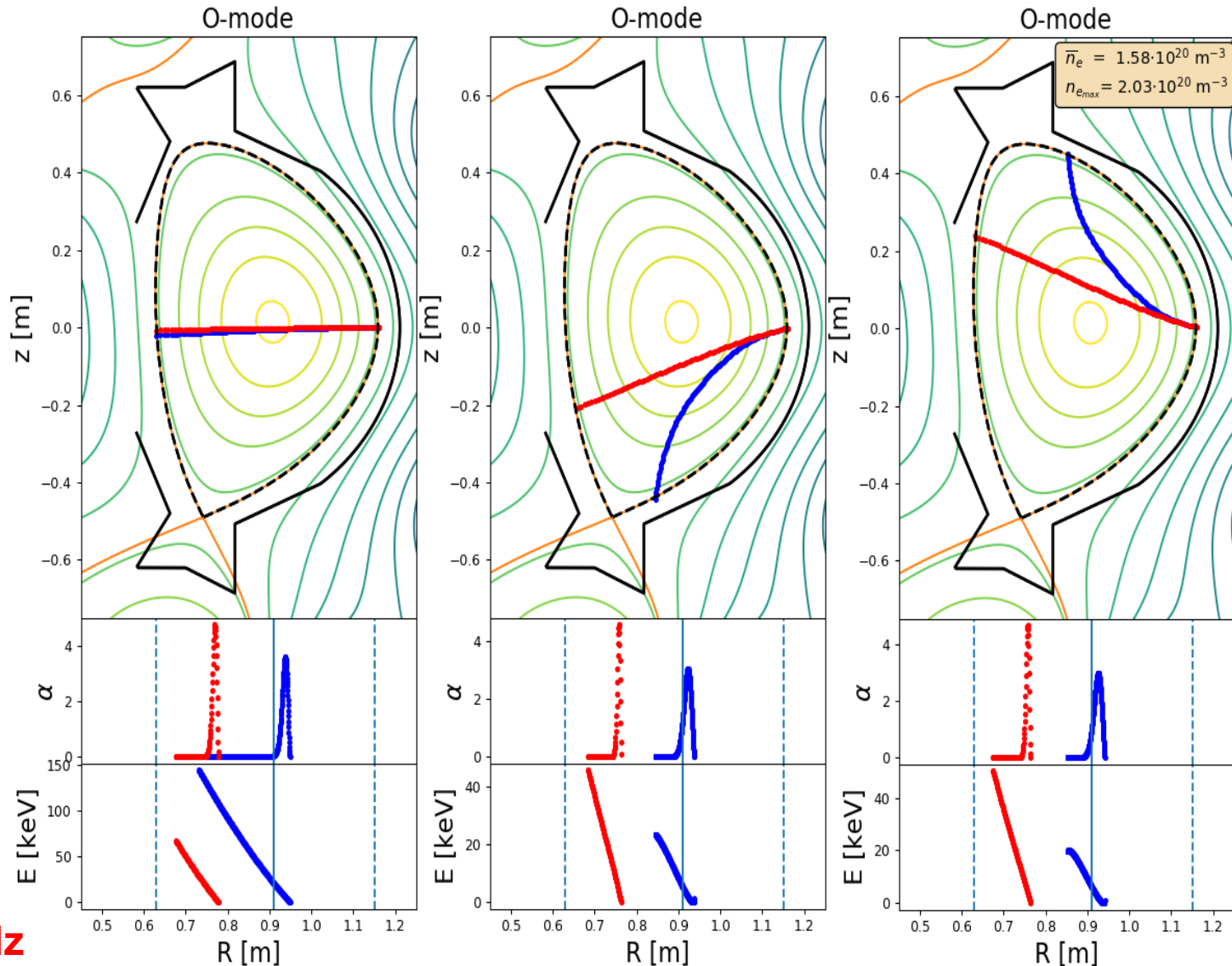
## Core density below cutoff condition



## Core density below cutoff condition

Everything just fine!

Perpendicular propagation poloidal angle  $\pm 14^\circ$



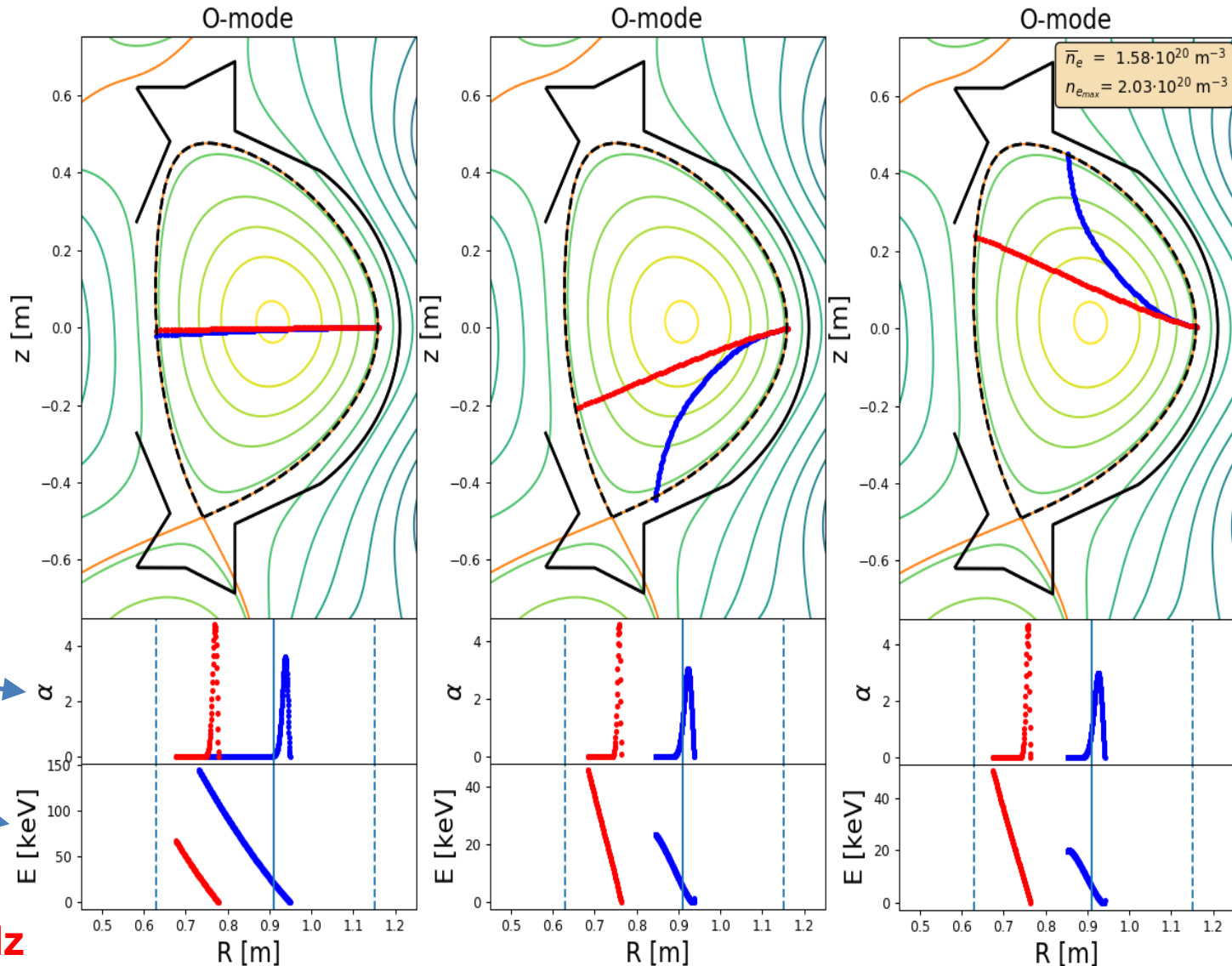
**Core density  
below cutoff  
condition**

**Everything just  
fine!**

Perpendicular  
propagation  
poloidal angle  
 $\pm 14^\circ$

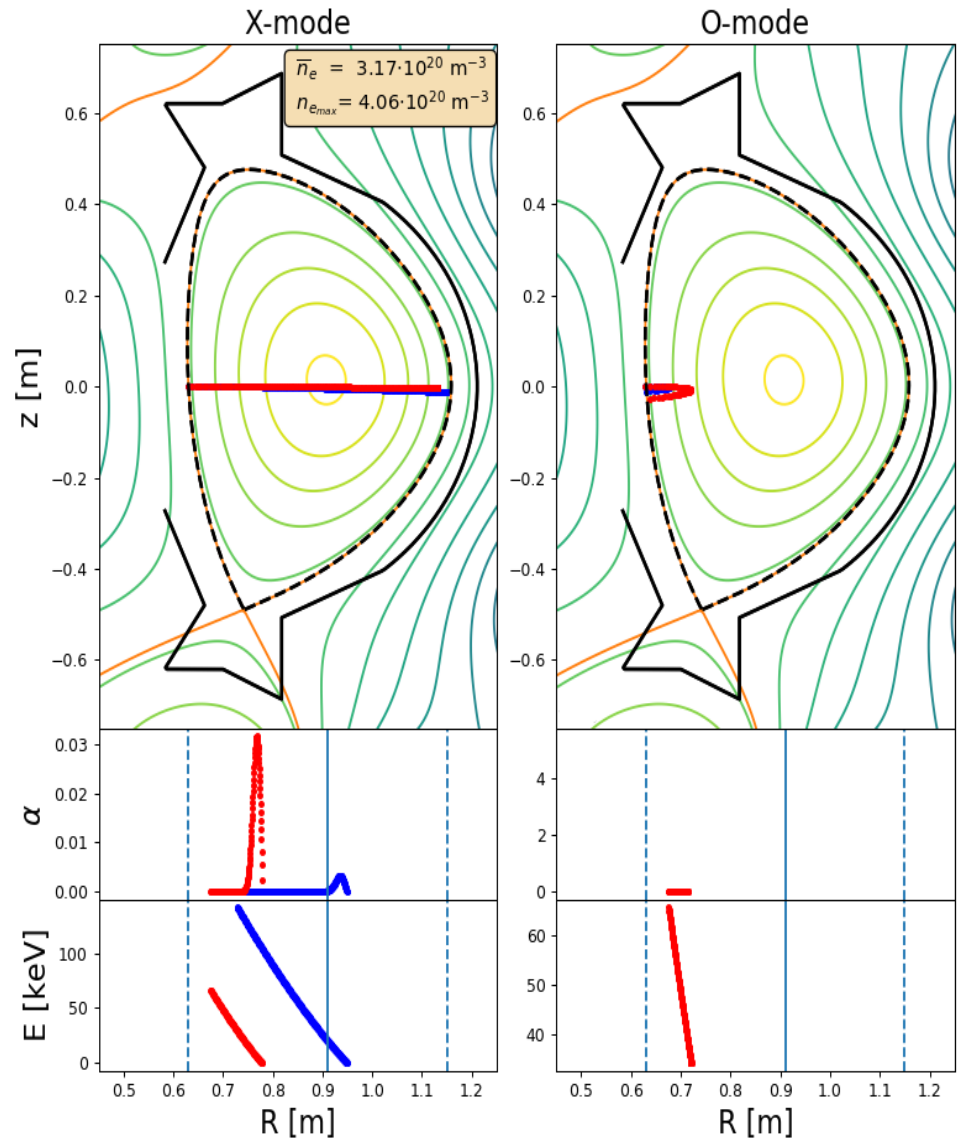
absorption coef.  $\alpha$

kinetic energy of  
resonant electrons E [keV]





## Core density above cutoff condition

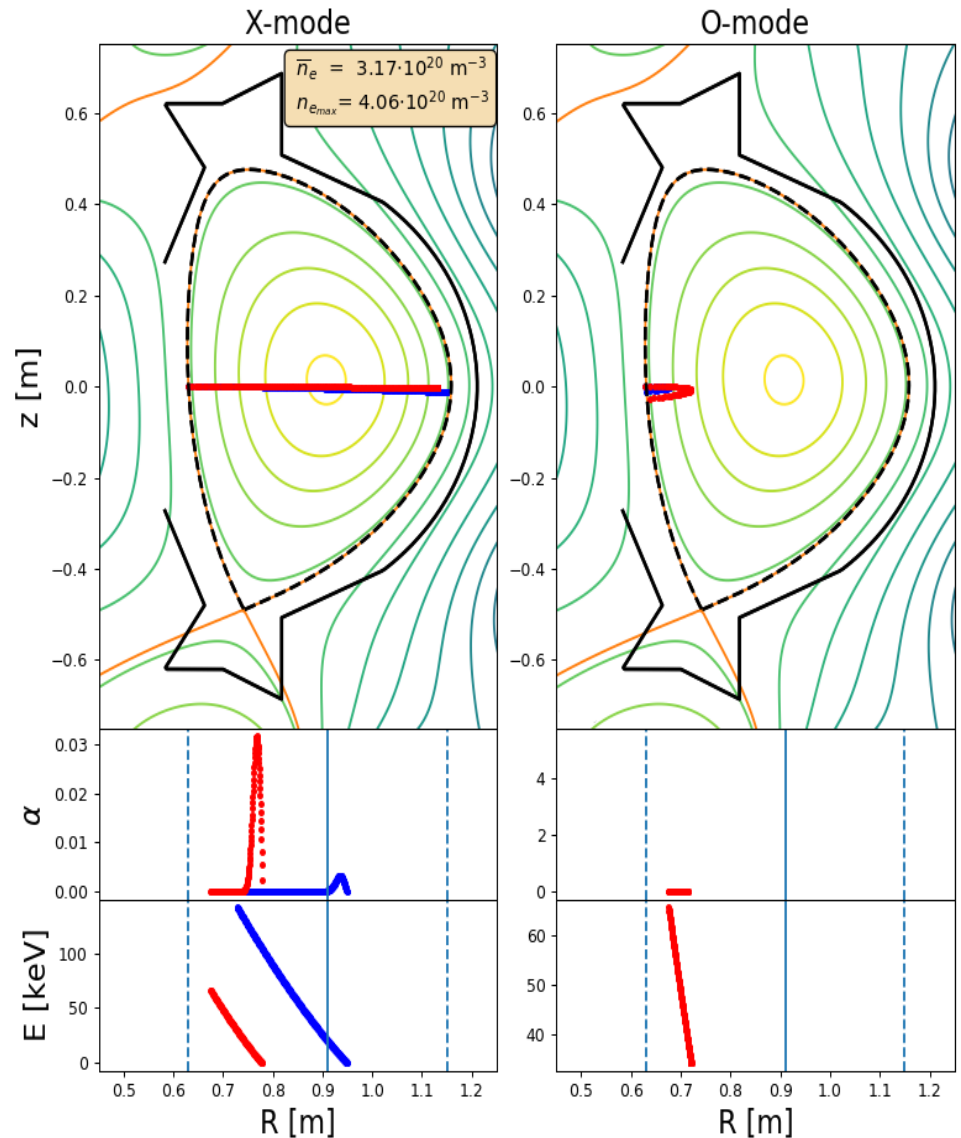


140 GHz

170 GHz

Core density above cutoff condition

Pretty bad!



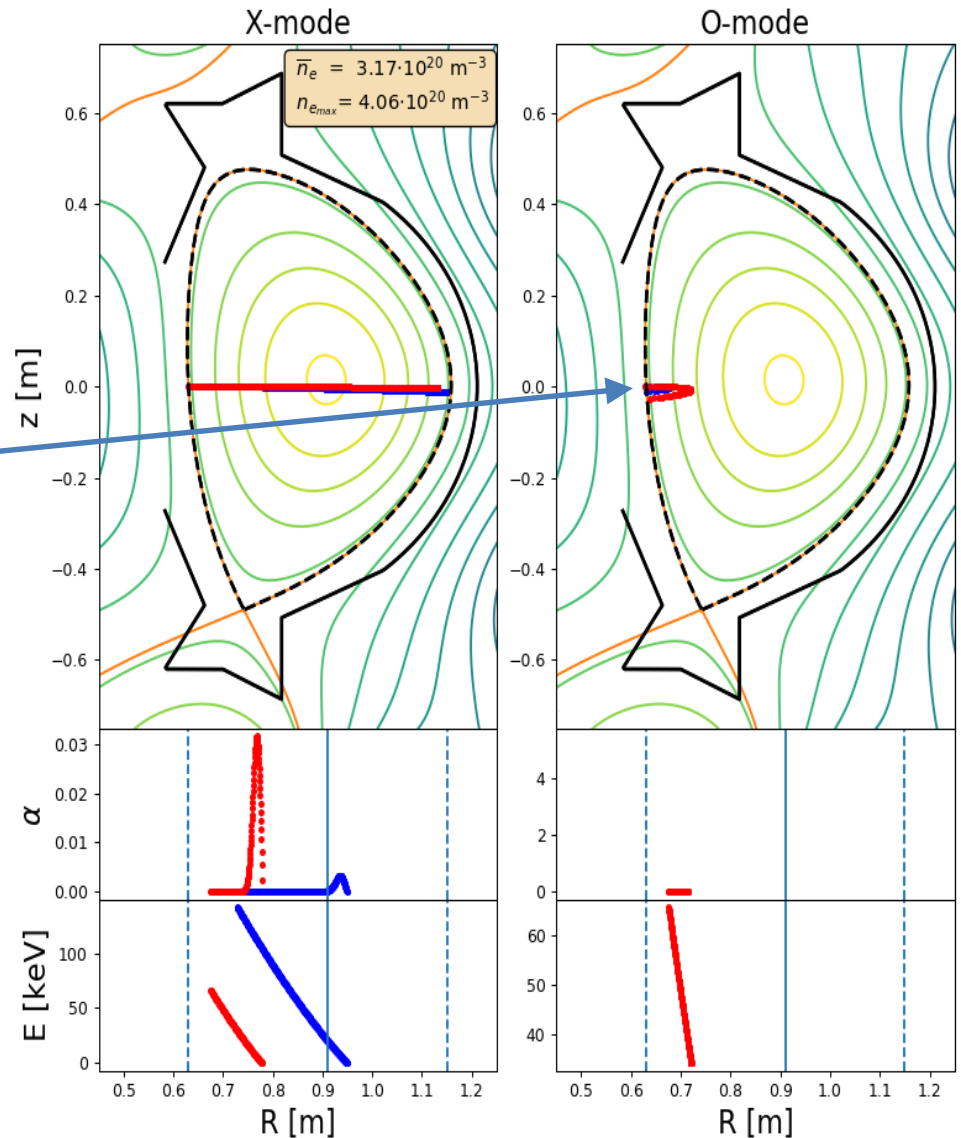
140 GHz

170 GHz

## Core density above cutoff condition

Pretty bad!

O-mode  $\rightarrow$  cutoff = reflection

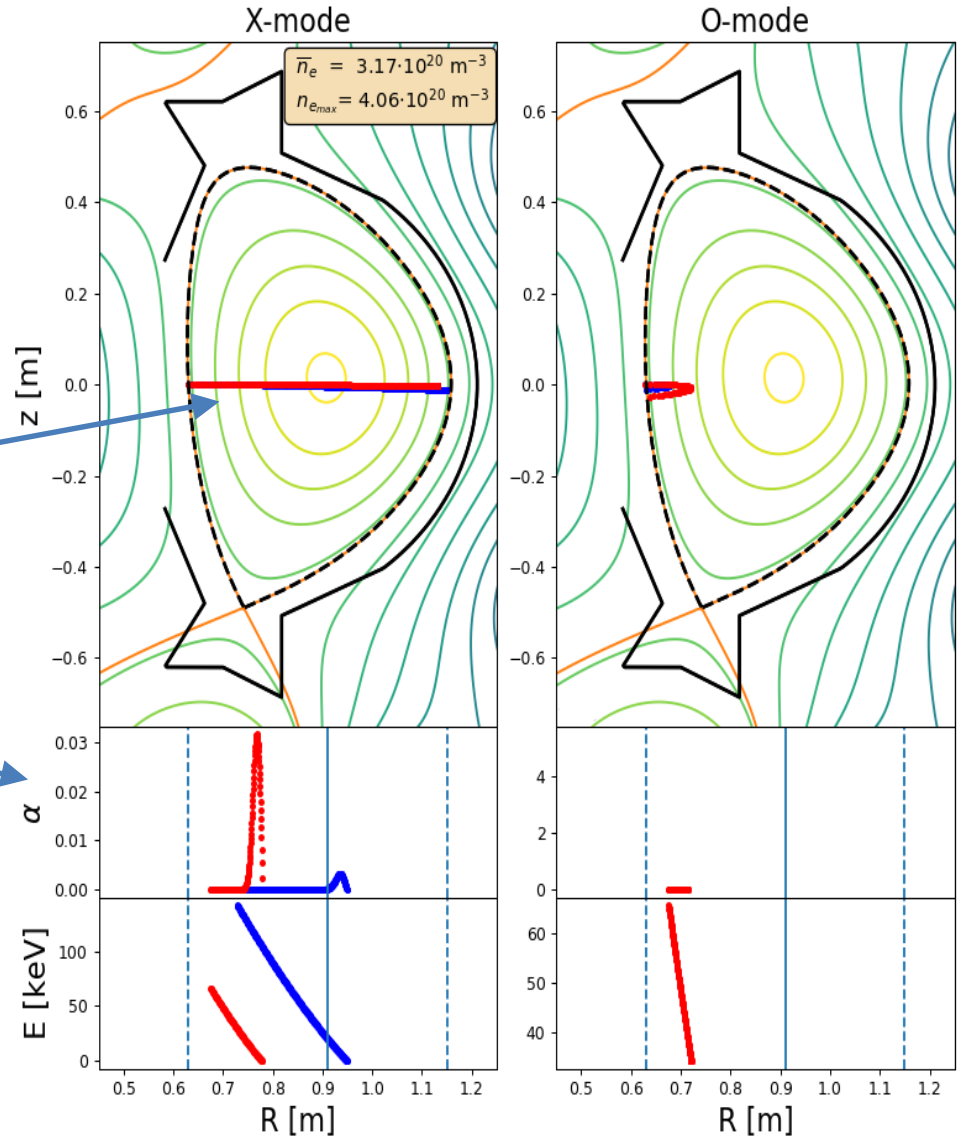


## Core density above cutoff condition

Pretty bad!

O-mode → cutoff = reflection

X-mode → from HFS but low  $\alpha$   
100x lower than O-mode



140 GHz      170 GHz

**Greenwald density:**  $n_G = \frac{I_p}{\pi a^2}$

**Natural H-mode density  $\approx 0.5 \cdot n_G$**

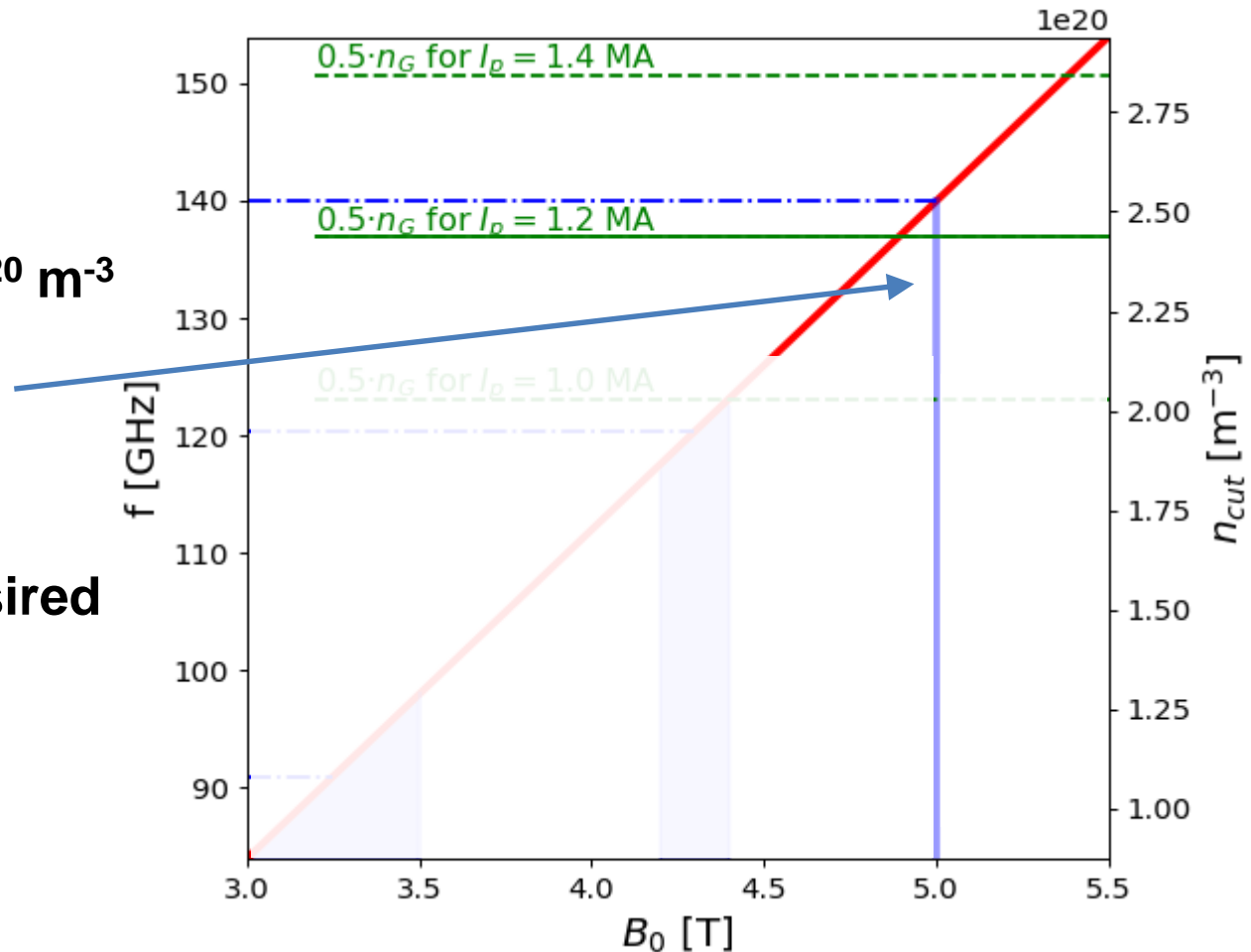
note: not maximum density!  
not for 300°C hot wall

**5 T  $\rightarrow$  140 Ghz**

$$n_{cut} = 2.5 \cdot 10^{20} \text{ m}^{-3}$$

**Only H-mode discharges with  $I_p \leq 1.2 \text{ MA}$**

**COMPASS-U: 2 MA desired**



## Proposal: ECRH for the High B-field (5 T) Discharges

Frequency: **140** – **150** GHz

O-mode cutoff: **2.5** - **2.8**·10<sup>20</sup> m<sup>-3</sup>

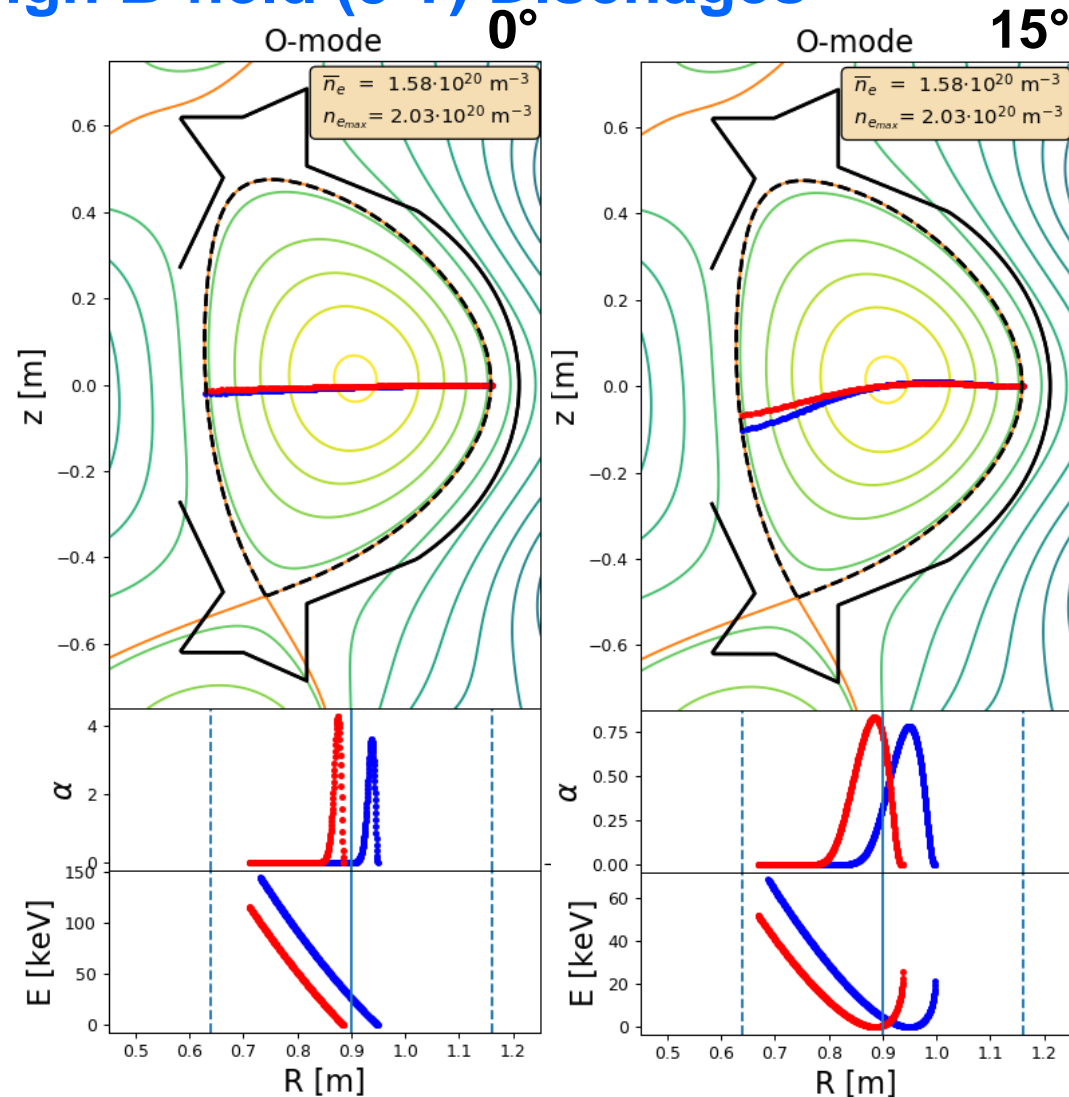
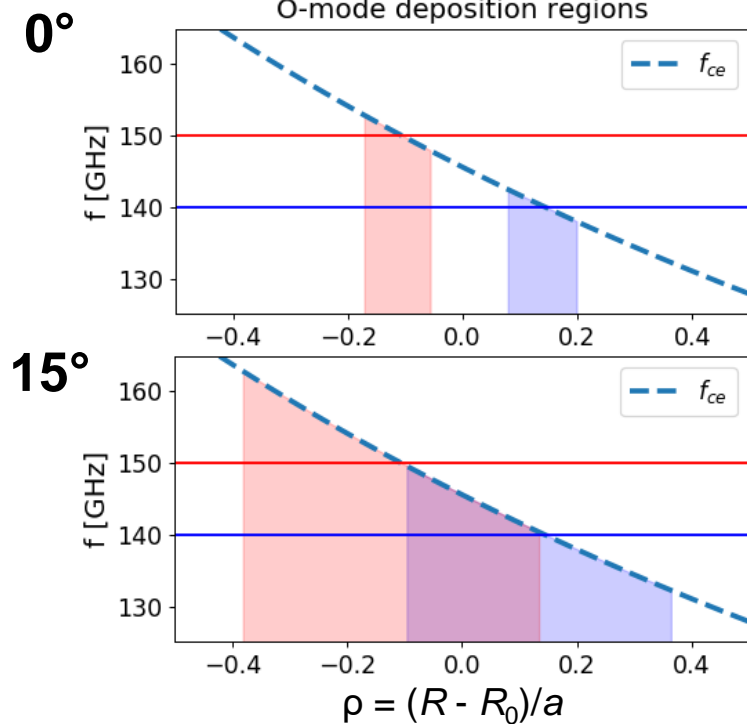
Toroidal angle: **0** – **15°**

## Proposal: ECRH for the High B-field (5 T) Discharges

Frequency: **140 – 150 GHz**

O-mode cutoff:  **$2.5 - 2.8 \cdot 10^{20} \text{ m}^{-3}$**

Toroidal angle: **0 – 15°**



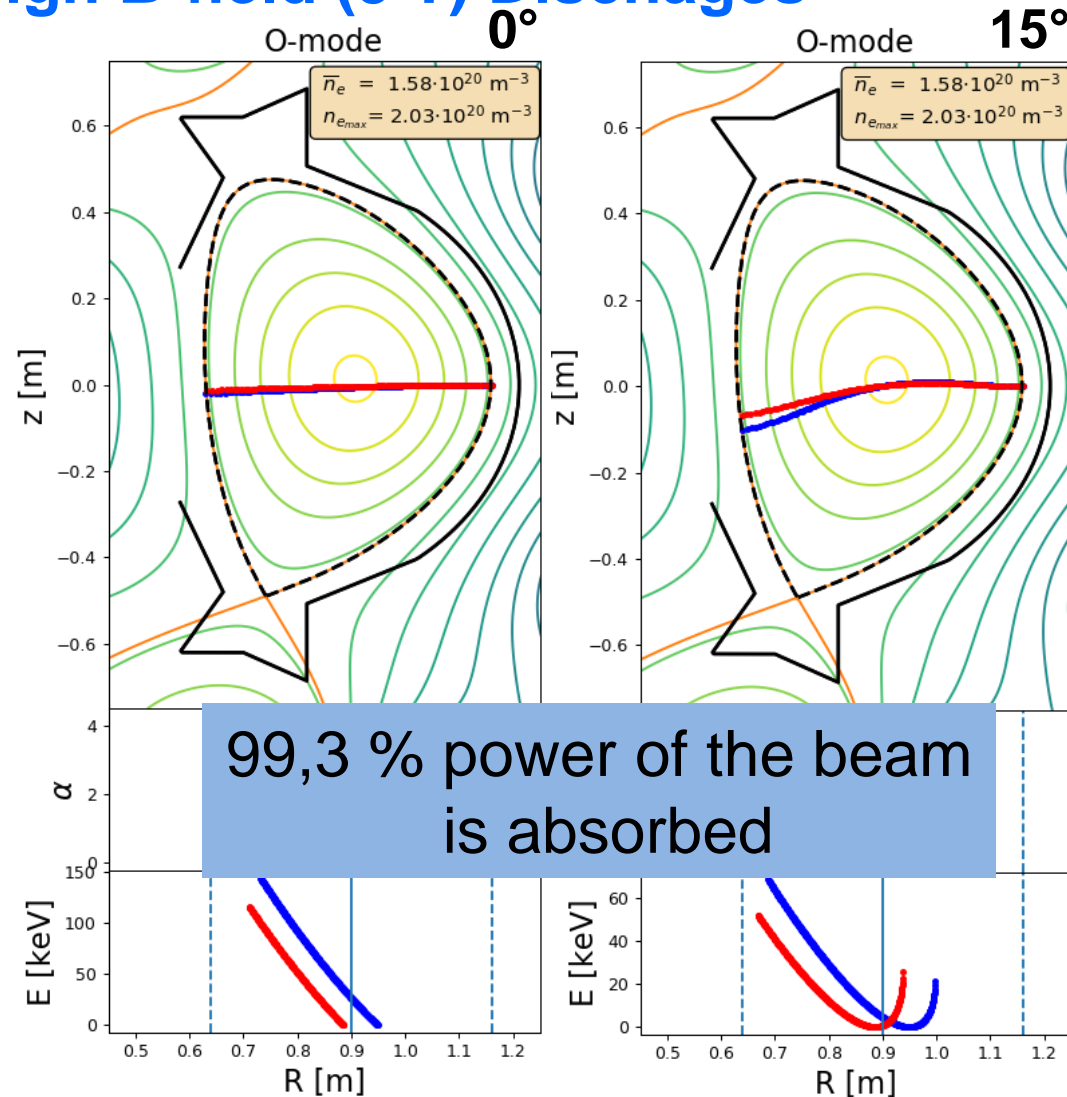
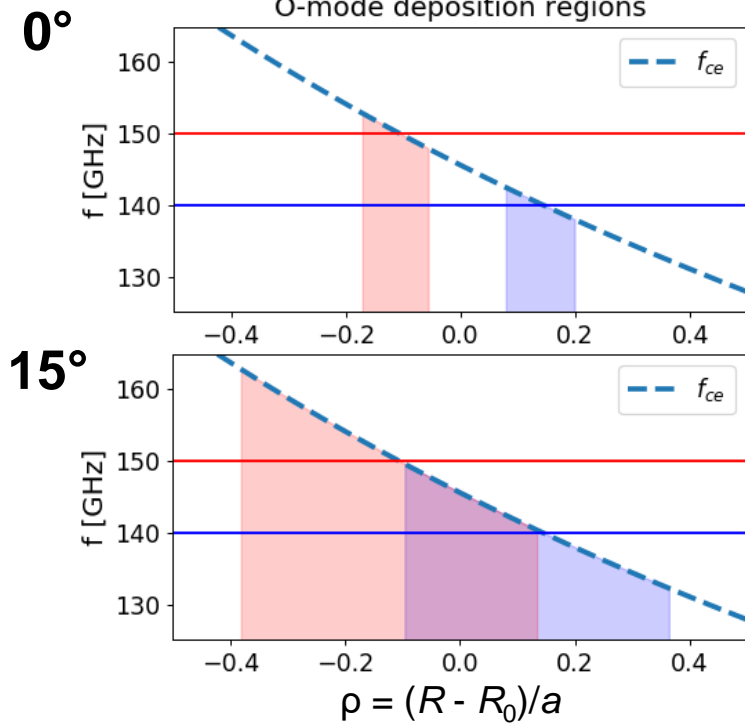


## Proposal: ECRH for the High B-field (5 T) Discharges

Frequency: **140 – 150 GHz**

O-mode cutoff:  **$2.5 - 2.8 \cdot 10^{20} \text{ m}^{-3}$**

Toroidal angle: **0 – 15°**





- **ECRH would be very favourable at COMPASS-U.**

- **ECRH would be very favourable at COMPASS-U.**
- **Reaching the core of the COMPASS-U plasma is strongly dependent on the electron density.**

- **ECRH would be very favourable at COMPASS-U.**
- **Reaching the core of the COMPASS-U plasma is strongly dependent on the electron density.**
- **Natural H-mode density in preferred discharge scenarios is much higher than the cutoff condition. ECRH will not be useful in later phases of these discharges.**

- **ECRH would be very favourable at COMPASS-U.**
- **Reaching the core of the COMPASS-U plasma is strongly dependent on the electron density.**
- **Natural H-mode density in preferred discharge scenarios is much higher than the cutoff condition. ECRH will not be useful in later phases of these discharges.**
- **It is necessary to model the behaviour of plasma in H-mode at COMPASS-U.**

- **ECRH would be very favourable at COMPASS-U.**
- **Reaching the core of the COMPASS-U plasma is strongly dependent on the electron density.**
- **Natural H-mode density in preferred discharge scenarios is much higher than the cutoff condition. ECRH will not be useful in later phases of these discharges.**
- **It is necessary to model the behaviour of plasma in H-mode at COMPASS-U.**
- **To avoid these restrictions, we should cooperate with the DEMO gyrotron development group in Russia.**

- [1] MAZZUCATO, E., *Electromagnetic Waves for Thermonuclear Fusion Research*. Singapore: World Scientific, 2014. ISBN 9789814571807
- [2] STIX, T. H., *Waves in Plasmas*. New York: American Institute of Physics, 1992.
- [3] BORNATICI, M., et al., Electron cyclotron emission and absorption in fusion plasmas. *Nuclear Fusion*. 1983, vyd. 23, c. 9, s. 1153.
- [4] FARINA, D., et al., SPECE: A code for electron cyclotron emission in tokamaks. In *AIP Conference Proceedings*, volume 988, pages 128–131, 2008.