

Plasma accretion near black holes

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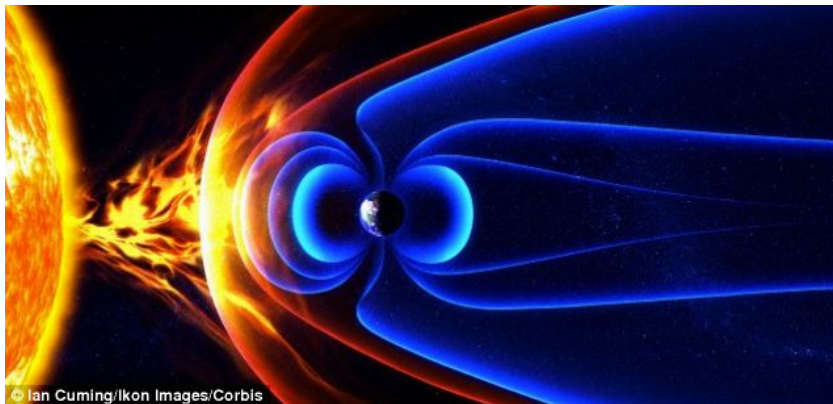


FuseNet PhD Event 2015, Prague

Outline

- 1 Plasma in the Universe
- 2 General relativistic MHD
- 3 Accretion disc

Magnetospheres of planets



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Credit: Ian Cuming

Sun



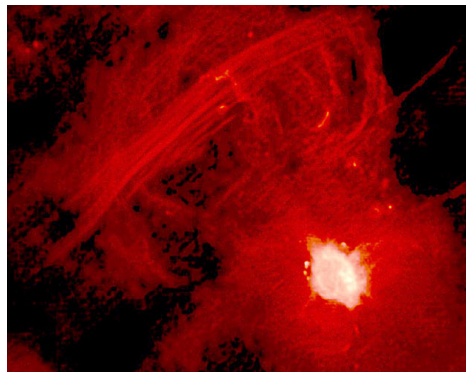
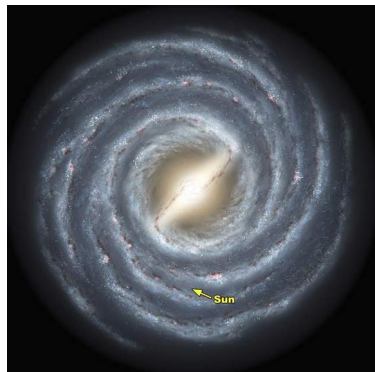
Sun

- G-type main-sequence star
- 99% of the mass of the Solar system
- Population I (heavy-element-rich star), Au, U
- Primarily from H and He, metals - less than 2%

Sun

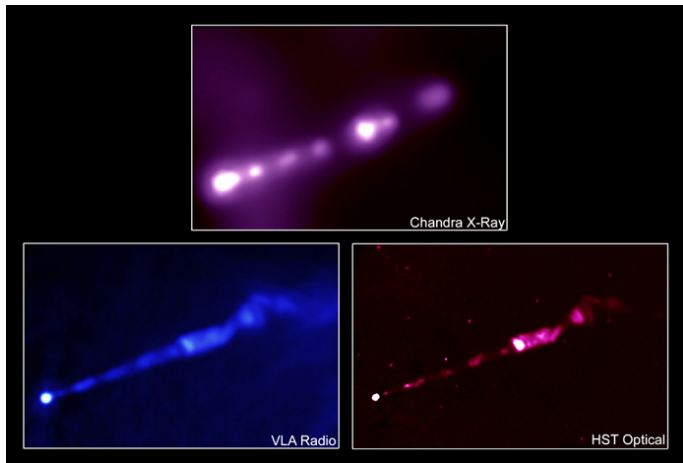
- Core
 - 20 – 25% of the Solar radius
 - Density up to $150 \text{ g} \cdot \text{cm}^{-3}$
 - Temperature $T \sim 15 \cdot 10^6 \text{ K}$
- Radiative zone, Tachocline, Convective zone
- Photosphere - visible surface of the Sun (5 800K)
- Chromosphere
- Corona - visible during a Solar eclipse, $T \sim 1 \cdot 10^6 - 6 \cdot 10^6 \text{ K}$

Galaxies



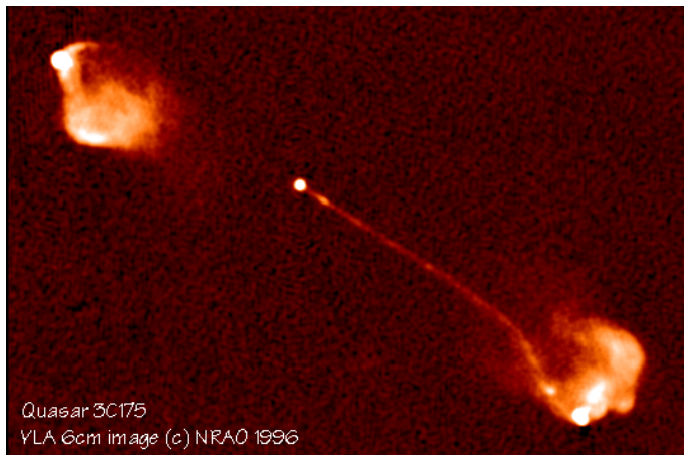
Credit: NASA/JPL-Caltech/R. Hurt (SSC/Caltech), F. Yusef-Zadeh et al., VLA, NRAO.

Galaxies



Jet from the radio galaxy Virgo A in different bands. Source:
<http://physweb.bgu.ac.il>

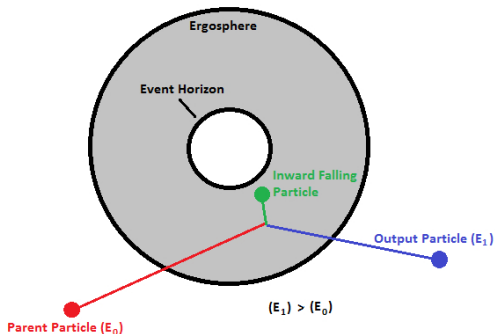
Galaxies



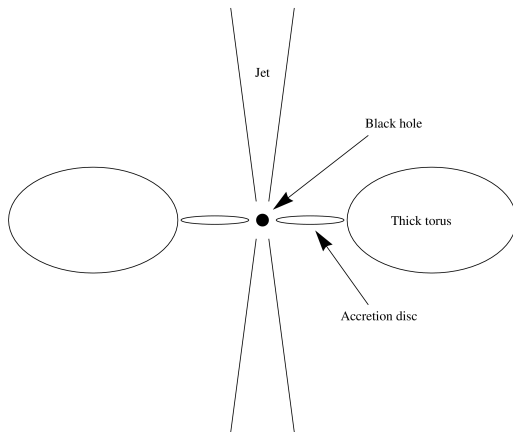
Jet from the quasar 3C175. Source: <http://physweb.bgu.ac.il>

Which mechanism drives jets?

- Blanford-Znajek mechanism and Penrose process

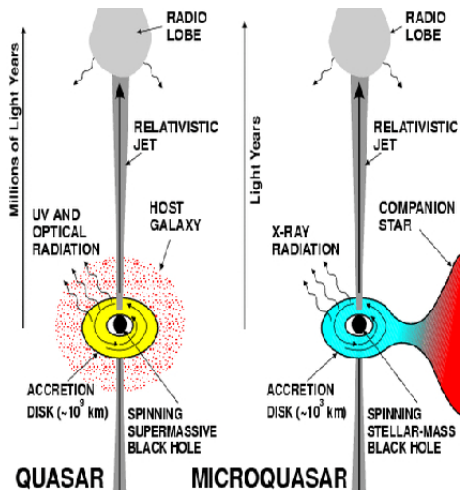


Active Galactic Nuclei



Schematic illustration of AGN

Quasar and microquasar



Quasars and other AGN

- Supermassive BH at centres of galaxies ($M_{BH} = 10^6 - 10^9 M_{\odot}$)
- Temperature in SMBH accretion discs: $T = 10^5 - 10^2 K$
- Accretion produces radiative power that often outshines the host galaxy
- Accretion disc is surrounded by moving gas clouds and large torus of gas and dust
- Very fast jets (almost speed of light) emerge from many AGN

Quasars and other AGN

- Black hole accretion in quasars is the most powerful and efficient stationary engine in the universe
- High angular momentum of rotating matter is transported outwards by stresses
- It allows matter to move inwards
- Gravitational energy is converted to heat
- Accretion disc physics: gravity, hydrodynamics, viscosity, radiation and magnetic fields

Quasars and other AGN

- More than 200 000 quasars are known
- Redshifts between 0.06 and 7.1 $\rightarrow 6 \cdot 10^8 - 29 \cdot 10^9$ light years away
- Luminosities are typically $10^{12}L_{\odot}$
- Typical luminosity: $L \sim 10^{40}$ *watts*

MHD equations

- Conductive fluids can support magnetic field
- Magnetic fields act on the fluid (plasma)
- Nonrelativistic MHD equations: conservation laws of mass, momentum and energy together with the induction equation

- $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0$

- $\frac{\partial(\rho \vec{v})}{\partial t} + \nabla \cdot (\rho \vec{v} \vec{v}) = -\nabla p + \vec{j} \times \vec{B} + \nabla \cdot \sigma$

- $\frac{\partial \rho}{\partial t} + \vec{v} \cdot \nabla \rho + \gamma p \nabla \cdot \vec{v} = Q$

- $\frac{\partial \vec{B}}{\partial t} = -\nabla \times \vec{E} = \nabla \times (\vec{v} \times \vec{B}) + \eta \nabla^2 \vec{B}$

MHD in general relativity

- Conservation of mass:

$$\partial_t(\sqrt{-g}\rho_0 u^t) = -\partial_i(\sqrt{-g}\rho_0 u^i)$$

$$\partial_t \rho = -\nabla \cdot (\rho \mathbf{v})$$

- Ideal MHD:

$$u_\mu F^{\mu\nu} = 0$$

$$\mathbf{E} + \mathbf{v} \times \mathbf{B}/c = 0$$

- Energy and momentum conservation:

$$\begin{aligned} \partial_t(\sqrt{-g}T^t_\nu) = \\ -\partial_i(\sqrt{-g}T^i_\nu) + \sqrt{-g}T^\kappa_\lambda \Gamma^\lambda_{\nu\kappa} \end{aligned}$$

$$\partial_t(\rho v) = -\nabla \cdot T - \rho \nabla \phi$$

- Induction equation:

$$\begin{aligned} \partial_t(\sqrt{-g}B^i) = \\ -\partial_j(\sqrt{-g}(u^j b^i - b^j u^i)) \end{aligned}$$

$$\partial_t \mathbf{B} = -\nabla \cdot (\mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v})$$

- No monopoles constraint:

$$\partial_i(\sqrt{-g}B^i) = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

Stationary solution

- The energy-momentum tensor of an ideal fluid:

- $$T^{\mu\nu} = (\rho_0 + p + u)u^\mu u^\nu + pg^{\mu\nu}$$

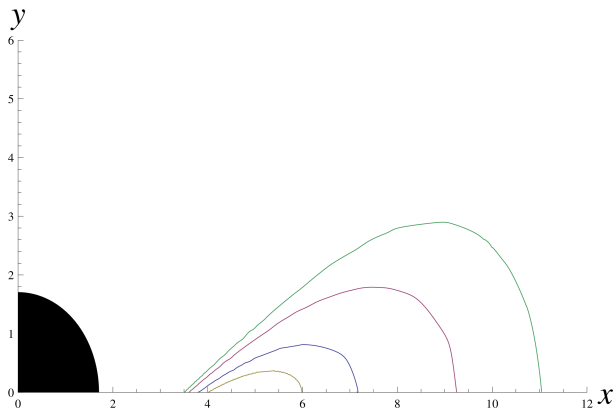
- Assuming a steady state and $\mathcal{L}_\xi u^\mu = 0$:

- $$T^{\mu\nu}_{;\nu} = 0 \rightarrow \frac{p_{,\mu}}{\rho+\epsilon} = (\ln u^t)_{,\mu} - \frac{l\Omega_{,\mu}}{1-\Omega l} - \text{relativistic Euler equation}$$

- Final equation for the torus:

- $$W - W_{\text{in}} \equiv \ln u_t - \ln u_{t\text{in}} - \int_{l_{\text{in}}}^l \frac{\Omega dl}{1-\Omega l} = - \int_0^p \frac{dp}{\rho+\epsilon}$$

Equipotential surfaces of the disc



Spin $a = 0.5$. The inner edge is at $R_{\text{in}} = 3.5, 3.6, 3.8, 4.0$.

Evolution equations

- Conservation laws

- $T^{\mu\nu} = (\rho_0 + u + p + b^2)u^\mu u^\nu + \left(p + \frac{b^2}{2}\right)g^{\mu\nu} - b^\mu b^\nu$

- $T^{\mu\nu}_{;\nu} = 0 \rightarrow 4$ equations

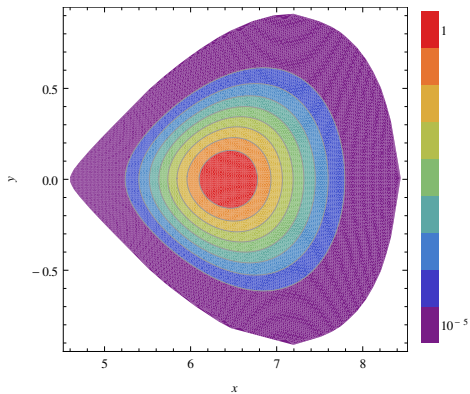
- $(\rho_0 u^\mu)_{;\mu} = 0 \rightarrow 1$ equation

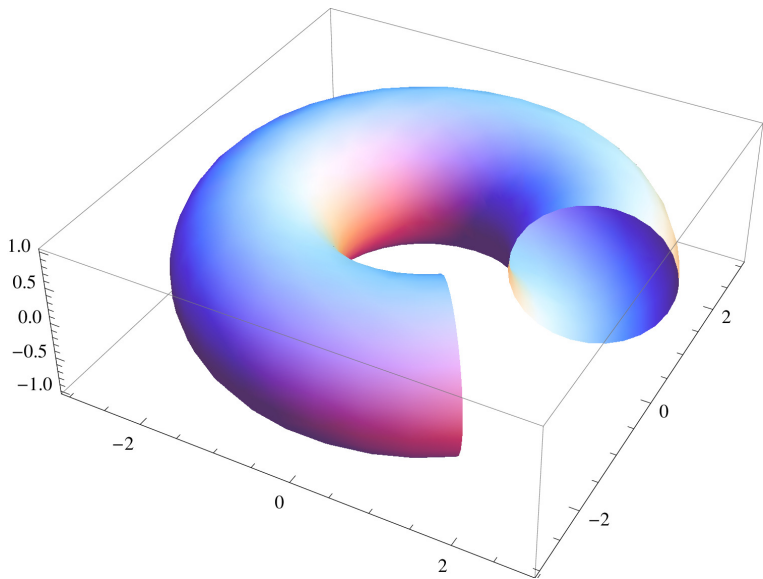
- $F^{\mu\nu}_{;\nu} = 0 \rightarrow 3$ evolution equations and constraint $\text{div } B = 0$

- Force-free constraint: $F^{\mu\nu} u_\nu = 0$

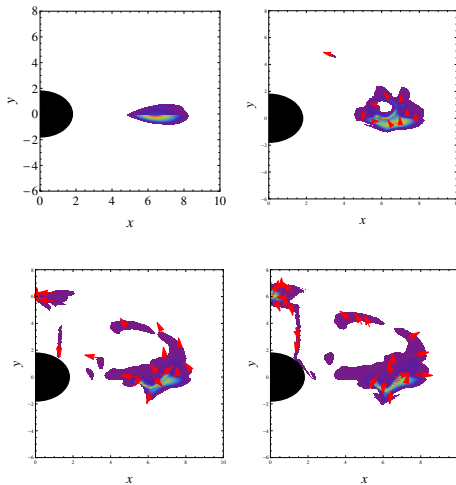
- Equation of state $p = \kappa \rho_0^\gamma$

Stationary distribution of mass



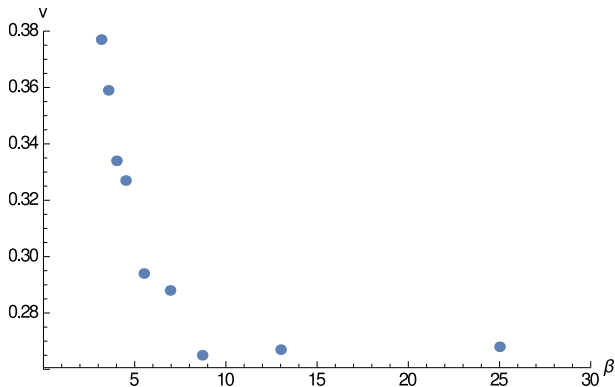


Mass accretion and ejection



$t = 0, 15, 19, 20, \beta = 5.5$

Dependence of outflow velocities on magnetization



Maximal value of vertical outflow velocity.

Thank you for your attention!

